





 Р. М.
 

[illegible]

h). “on the other hand, classical statistical mechanics is essentially only concerned with Type b) [i.e. non integrable systems], for in this case the microcanonical average is the same as the time average”.

(1) 
$$i\hbar \frac{\partial}{\partial t} \psi = \left( -\frac{\hbar^2}{2m} \Delta + V \right) \psi$$
 where  $\Delta$  is the Laplacian operator,  $\psi$  is the wave function, and  $V$  is the potential energy. The energy eigenvalues  $E_n$  are determined by the time-independent Schrödinger equation:
 
$$H \psi_n = E_n \psi_n$$
 where  $H = -\frac{\hbar^2}{2m} \Delta + V$  is the Hamiltonian operator. The energy levels  $E_n$  are real numbers,  $E_n \in \mathbb{R}$ .

[illegible]

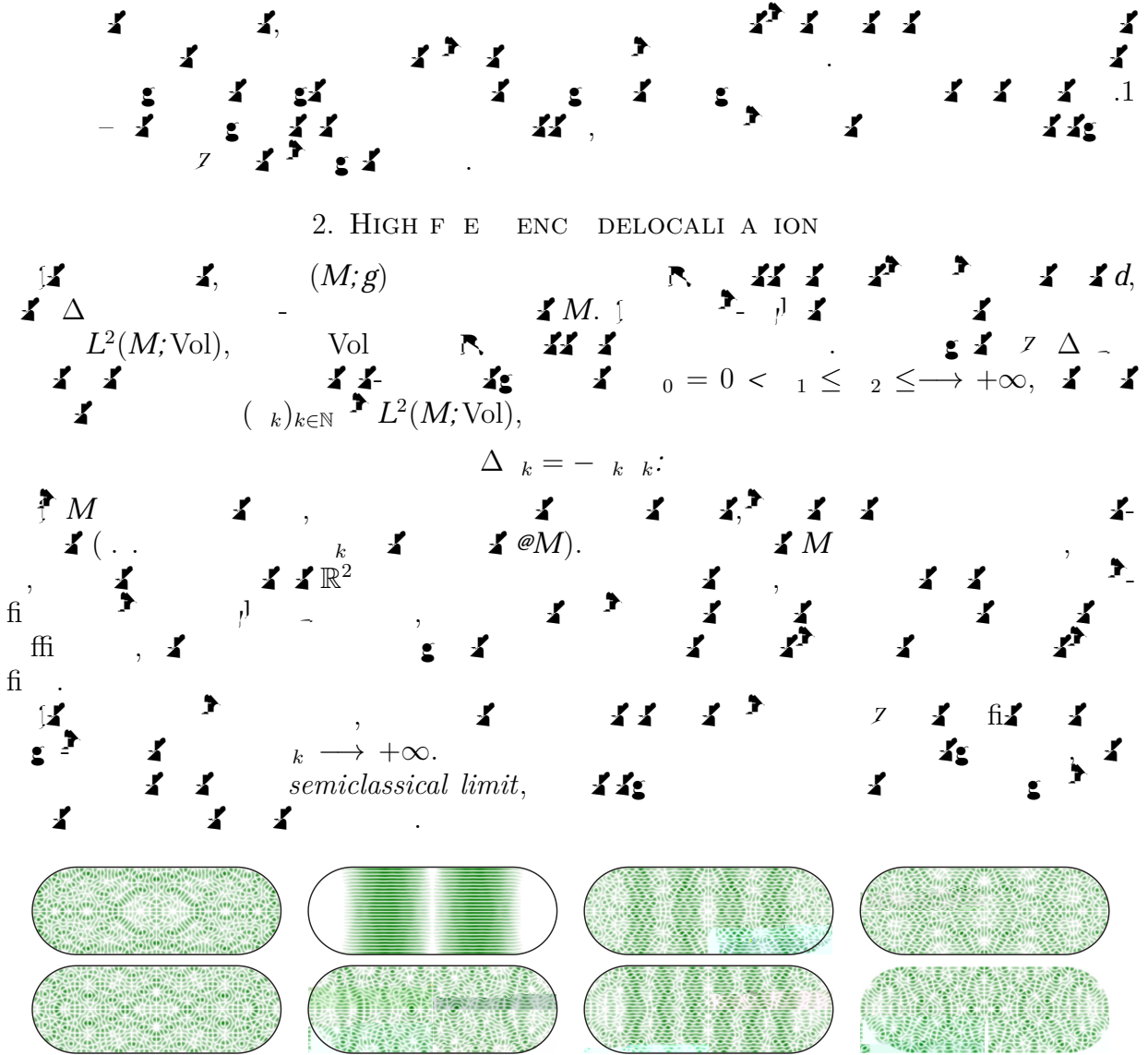


FIGURE 1.  $|n(x; y)|^2$  for  $n = 319$ . The plots show the distribution of eigenvalues for  $n = 319$ .

## 2.1. The role of the geodesic flow .

$$-\hbar^2 \Delta_k = E_k \quad (1) \quad \left( \begin{array}{c} \text{where } \Delta_k = -\hbar^2 \Delta_k \end{array} \right)$$

“ ”.

$100$ .

$-\hbar^2\Delta$

$M$

$T^*M$  *cotangent bundle*

$(x; \cdot) \in T^*M$   $x \in M$

$(x; \cdot) \in T^*M$

$t \in \mathbb{R}$ ,  $g^t(x; \cdot) \in T^*M$

$(g^t)_{t \in \mathbb{R}} : T^*M \longrightarrow T^*M$   $g^{t+s} = g^t \circ g^s$   $g^0$

*unit cotangent bundle*  $S^*M = \{(x; \cdot) \in T^*M; \| \cdot \|_x = 1\}$

$g^t$ .

$(k \longrightarrow +\infty)$ ,

$\frac{\partial}{\partial t} = i\Delta$

*wavefronts*

$\epsilon(x$

2.2.  $L^p$ -norms as measures of delocalization ?

$$\|k\|_\infty = O(k^{(d-1)/4}):$$
$$\|\lambda\|_{L^p} = O\left(\frac{\mu(\mathcal{I})}{2}\right)$$

- $(p) = d\left(\frac{1}{2} - \frac{1}{p}\right) - \frac{1}{2}$  for  $\frac{2(d+1)}{(d-1)} \leq p \leq +\infty$ ;
- $(p) = \frac{d-1}{2} \left(\frac{1}{2} - \frac{1}{p}\right)$  for  $2 \leq p \leq \frac{2(d+1)}{(d-1)}$ .

**Theorem 3** ( [10] - [11] ). *Assume there exists a subsequence  $n_k \rightarrow +\infty$  and  $C > 0$  such that  $\|\lambda_{n_k}\|_\infty \geq C^{\frac{\mu(\infty)}{2}}$ . Then there exists  $x$  such that  $x(\mathcal{L}_x) > 0$ ;*

- If  $M$  is real analytic, the existence of such subsequence  $\lambda_{n_k}$  is sufficient to the existence of  $x$  such that  $\mathcal{L}_x = S_x^* M$ , and the first return map  $\tau_x : S_x^* M \rightarrow S_x^* M$  possesses an absolutely continuous invariant probability measure. (Moreover, in*

that case, there exists  $t_0 > 0$  such that  $g^{t_0}(x; v) \in S_x^*M$  for all  $v \in S_x^*M$ , that is, there is a common return time).

- If  $M$  is real analytic and  $\dim M = 2$ , the existence of such subsequence  $\lambda_{n_k}$  is due to the existence of  $x \in M$  and  $t_0 > 0$  such that  $g^{t_0}(x; v) = (x; v)$  for all  $v \in S_x^*M$ .

Anosov property,  $\| \lambda \|_{L^\infty} = o\left(\frac{\mu(\infty)}{2}\right) (O(1) o)$ .

**Theorem 4.** (i) If  $d = 2$  and  $M$  has no conjugate points, or if  $d \geq 2$  and  $M$  has non-positive sectional curvature, for  $p = +\infty$ ,

$$\| \lambda \|_{L^p} = O\left(\frac{\mu(p)}{\sqrt{\log}}\right):$$

(i') Statement (i) actually holds if  $M$  has no conjugate points, for all  $d \geq 2$ .

(ii) (i) holds for all  $p > p_c$ .

(iii) If  $M$  has non-positive sectional curvature, for  $p < p_c$ , there exists  $(p; d) > 0$  such that

$$\| \lambda \|_{L^p} = O\left(\frac{\mu(p)}{(\log)^{\sigma(p,d)}}\right)$$

(iv) Statement (iii) still holds for  $p = p_c$ .

( ) ( ) 1 “ ” fl .

### 2.3. The Shnirelman theorem and the Quantum Unique ergodicity conjecture .

$| \psi_k(x) |^2 d\text{Vol}(x)$ .  
 $k \rightarrow +\infty$  ,  $| \psi_k(x) |^2 d\text{Vol}(x)$

Let  $(x_0; 0) \in S^*M$ ,  $a : S^*M \rightarrow \mathbb{R}$ ,  $T \rightarrow +\infty$

$$\left( \frac{1}{T} \int_0^T a \circ g^t(x_0; 0) dt - \int_{S^*M} a dL \right) \xrightarrow{T \rightarrow +\infty} 0$$

for all  $(x_0; 0) \in S^*M$ .

### Quantum Ergodicity Theorem (Shnirelman theorem).

**Theorem 5** ([10], [11]). Let  $(M; g)$  be a compact Riemannian manifold, with the metric normalized so that  $\text{Vol}(M) = 1$ . Call  $\Delta$  the Laplace-Beltrami operator on  $M$ . Assume that the geodesic flow of  $M$  is ergodic with respect to the Liouville measure. Let  $(e_k)_{k \in \mathbb{N}}$  be an orthonormal basis of  $L^2(M; g)$  made of eigenfunctions of the Laplacian

$$\Delta e_k = -\lambda_k e_k; \quad \lambda_k \leq \lambda_{k+1} \rightarrow +\infty.$$

Let  $a$  be a continuous function on  $M$ . Then

$$(3) \quad \frac{1}{N(\lambda)} \sum_{k, \lambda_k \leq \lambda} \left| \langle e_k, a \rangle_{L^2(M)} - \int_M a(x) d\text{Vol}(x) \right|^2 \xrightarrow{\lambda \rightarrow +\infty} 0$$

where the normalizing factor is  $N(\lambda) = |\{k; \lambda_k \leq \lambda\}|$ :

$$\langle e_k, a \rangle_{L^2(M)} = \int_M a(x) |e_k(x)|^2 d\text{Vol}(x).$$

**Remark 6.** The Cesaro limit (3) implies that there exists a subset  $S \subset \mathbb{N}$  of density 1 such that

$$(4) \quad \langle e_k, a \rangle_{L^2(M)} \xrightarrow{n \rightarrow +\infty, n \in S} \int_M a(x) d\text{Vol}(x):$$

In addition, using the fact that the space of continuous functions is separable, one can actually find  $S \subset \mathbb{N}$  of density 1 such that (4) holds for all  $a \in C^0(M)$ . In other words, the sequence of measures  $(|e_k(x)|^2 d\text{Vol}(x))_{n \in S}$  converges weakly to the uniform measure  $d\text{Vol}(x)$ .

Actually, the full statement of the theorem says that there exists a subset  $S \subset \mathbb{N}$  of density 1 such that

$$(5) \quad \langle e_k, A e_k \rangle_{L^2(M)} \xrightarrow{n \rightarrow +\infty, n \in S} \int_{S^*M} \sigma^0(A) dL$$

for every pseudodifferential operator  $A$  of order 0 on  $M$ . On the right-hand side,  $\sigma^0(A)$  is the principal symbol of  $A$ , that is a function on the unit cotangent bundle  $S^*M$ . Equation 4 corresponds to the case where  $A$  is the operator of multiplication by the function  $a$ .



“ $\mu$ ” ( , ).  
 $\hbar \rightarrow 0$  ,  
 $2, 1, 2$  .  
 $k \rightarrow +\infty$  ,  
 $1, 3$  .

### Quantum Unique Ergodicity conjecture.

$M$  ,  
 $200$  ,  
 $M$  ,  
 $curvatures$  ,  
 $M$  ,  
 $\lambda$  ,  
 $111$  ,  
 $20$  .

### 2. . Entropy and support of semiclassical measures.

$d$  .  
 $\langle n; A_n \rangle$  ,  
 $\langle n_k; A_{n_k} \rangle$  ,  
 $\int_{S^*M} \rho(A) d\mu$  ,  
 $S^*M$  .

## 2.1.

[illegible]

[illegible]

$L$

[illegible]





[illegible]

### 3. LA GEOMETRIA DELLE LOCALITÀ A ION

The figure consists of two panels, labeled (a) and (b), showing numerical simulations of particle dynamics in a potential well.

Panel (a) displays several trajectories starting from different initial conditions. The trajectories are shown as sequences of points. Some trajectories appear regular and confined to specific regions, while others show more complex, possibly chaotic behavior. The label "dynamics" is present near the bottom left.

Panel (b) shows a similar setup but with a focus on the distribution of points. A large number of points are plotted, forming distinct patterns or clusters. The label "spectral statistics" is visible near the top right. Below it, there is a mention of "( / )".

Both panels include mathematical symbols such as  $\Delta = \frac{d^2}{dx^2}$ ,  $\rightarrow +\infty$ , and various Greek letters like  $\alpha$  and  $\beta$ .

3.  $\dots$  regular  $\dots$

### 3.1. Overview of the problem.

$G = (V; E)$ .  
localized, delocalized?

$I \subset \mathbb{R}$ ,

- spectral localization :
- exponential localization :
- dynamical localization :

- spectral delocalization :
- ballistic transport :

$(G_N)$   
 $(G_N)$   
 $N \rightarrow \infty$ .

$\sum_{x=1}^N | \dots |^2$   
 $M_N$   
 $N \times N$ ,  
 $( \dots )_{j=1}^N$

- $\infty$  norms :

$\| \dots \|_\infty$

- $p$  norms:

non-ergodic ( multi-fractal )

- Scarring :

Quantum ergodicity :

most  $j$ , all  $j$ , quantum unique ergodicity.  
 almost sure random random  
 2, 3, 33, 1, 1, 13, 1  
 3, ff

### 3.2. Entropy.

$(q+1)$ -  
 $(q+1)$ -  
 $(G_N)_{N \in \mathbb{N}} = (V_N; E_N)$ .

$$\mathcal{A}_N f(x) = \sum_{x \sim y} f(y)$$

$$x \sim y \iff x \text{ and } y$$

$$\Delta_N f(x) = \sum_{x \sim y} (f(y) - f(x)) :$$

$$(10) \quad \mathcal{A}_N - (q+1)I = \Delta_N :$$

**Theorem 10** ( ). Let  $(G_N)$  be a sequence of  $(q+1)$ -regular graphs (with  $q$  fixed),  $G_N = (V_N; E_N)$  with  $V_N = \{1; \dots; N\}$ . Assume that<sup>1</sup> there exists  $c > 0$ ;  $\delta > 0$  such that, for any  $k \leq c \ln N$ , for any pair of vertices  $x, y \in V_N$ ,

$$(11) \quad |\{\text{paths of length } k \text{ in } G_N \text{ from } x \text{ to } y\}| \leq q^{k(\frac{1-\delta}{2})} :$$

Fix  $\delta > 0$ . Then, if  $f$  is an eigenfunction of the discrete Laplacian on  $G_N$  and if  $\Lambda \subset V_N$  is a set such that

$$\sum_{x \in \Lambda} |f(x)|^2 \geq \sum_{x \in V_N} |f(x)|^2 ;$$

then  $|\Lambda| \geq N^\alpha$  — where  $\alpha > 0$  is given as an explicit function of  $\delta$ ; and  $c$ .

$$H_N(f) = -\frac{1}{\log N} \sum_x |f(x)|^2 \ln |f(x)|^2$$

<sup>1</sup>  $c \ln N \geq c \ln N$

$$\|L^\infty\|_{\infty} = O((\log N)^{-1/4}). \quad (11)$$

3.3. QE on regular graphs. Let  $G_N = (V_N; E_N)$  be a  $(q+1)$ -regular graph with  $N$  vertices and  $E_N$  edges.

$$(EXP) \quad \frac{1}{N} \sum_{x \in V_N} \left( \frac{1}{(q+1)^{-1} \mathcal{A}_N} \right)^2 (V_N) \rightarrow \{1\} \cup [-1+\epsilon; 1-\epsilon]$$

$$0, \quad \frac{1}{N} \sum_{x \in V_N} \left( \frac{1}{\mathcal{A}_N} \right)^2 (V_N) \rightarrow \{1\} \cup [-1+\epsilon; 1-\epsilon]$$

(BST)  $R$ ,

$$\frac{|\{x \in V_N; (x) < R\}|}{N} \xrightarrow{N \rightarrow \infty} 0$$

$$\frac{(x)}{B(x; r)} \rightarrow \frac{1}{r} \quad \text{as } r \rightarrow \infty$$

$$(BST) \quad \frac{1}{N} \sum_{x \in V_N} \left( \frac{1}{(q+1)^{-1} \mathcal{A}_N} \right)^2 (V_N) \rightarrow \{1\} \cup [-1+\epsilon; 1-\epsilon]$$

$$\frac{1}{N} |\{j; \binom{(N)}{j} \in I\}| \xrightarrow{N \rightarrow +\infty} \int_I m(\cdot) d\mu$$

$$(12) \quad \int \langle \cdot, \mathcal{A}_x^k \cdot \rangle_{\ell^2(x)} d\mu = \langle \cdot, \mathcal{A}_x^k \cdot \rangle_{\ell^2(x)}$$

$$\mathcal{A}_x \rightarrow \mathcal{A}_x^k \quad \text{as } k \rightarrow \infty$$



**Theorem 11** (). Let  $(G_N) = (V_N; E_N)$  be a sequence of  $(q+1)$ -regular graphs with  $|V_N| = N$ . Assume that  $(G_N)$  satisfies **(BST)** and **(EXP)**.

Let  $(\binom{N}{1}, \dots, \binom{N}{N})$  be an orthonormal basis of eigenfunctions of  $\mathcal{A}_N$  in  $\mathcal{L}^2(V_N)$ .

Let  $\mathbf{a}_N : V_N \rightarrow \mathbb{C}$  be a sequence of functions such that  $\sup_N \sup_{x \in V_N} |\mathbf{a}_N(\mathbf{x})| \leq 1$ : Define  $\langle \mathbf{a}_N \rangle = \frac{1}{N} \sum_{x \in V_N} \mathbf{a}_N(\mathbf{x})$ .

*Then*

$$\begin{aligned}
& \frac{1}{N} \sum_{j=1}^N \left| \langle \psi_j^{(N)}; a_N \psi_j^{(N)} \rangle_{\ell^2(V_N)} - \langle a_N \rangle \right|^2 \xrightarrow{N \rightarrow +\infty} 0; \\
(13) \quad & \frac{1}{N} \left| \left\{ j \in [1; N]; \left| \langle \psi_j^{(N)}; a_N \psi_j^{(N)} \rangle_{\ell^2(V_N)} - \langle a_N \rangle \right| > \right\} \right| \xrightarrow{N \rightarrow +\infty} 0; \\
& \langle \psi_j^{(N)}; a_N \psi_j^{(N)} \rangle_{\ell^2(V_N)} = \sum_{x \in V_N} a_N(x) |\psi_j^{(N)}(x)|^2. \\
& \sum_{x \in V_N} |\psi_j^{(N)}(x)|^2 = \frac{1}{N} \sum_{x \in V_N} 1 = 1.
\end{aligned}$$

### 3. . Non-regular graphs : from spectral to spatial delocalization.

regular  $\vdots$   
 quantum ergodicity theorem  $\vdots$   
 spectral delocalization  $\vdots$   
 spatial delocalization  $\vdots$

*“If a large finite system is close (in the Benjamini-Schramm topology) to an infinite system having purely absolutely continuous spectrum in an interval  $I$ , then the eigenfunctions (with eigenvalues lying in  $I$ ) of the finite system satisfy quantum ergodicity.”*

$$(G_N)_{N \in \mathbb{N}}, \quad V_N, E_N, \quad N \longrightarrow \infty, \quad \text{fix } \rho, \quad G_N, \quad d, \quad \mathcal{A}_N : \mathbb{C}^{V_N} \rightarrow \mathbb{C}^{V_N} \quad |V_N| =$$

$$(\mathcal{A}_N f)(v) = \sum_{w \sim v} f(w);$$



$$\begin{aligned}
& \mathcal{T} \ni v; w \in \mathcal{T} \quad v \sim w, \quad \mathcal{T}^{(v|w)} \quad \mathcal{A}_{\mathcal{T}^{(v|w)}} \\
& \quad \gamma_w(v) := -\mathcal{G}^{(v|w)}(v; v; \cdot). \quad \mathcal{G}^{(v|w)}(\cdot; \cdot; \cdot) \\
& \text{(Green)} \quad I, \quad s > 0 \\
& \sup_{\lambda \in I, \eta_0 \in (0,1)} \mathbb{E} \left( \sum_{y: y \sim o} |\text{Im } \mathcal{G}_o^{\lambda + i\eta_0}(y)|^{-s} \right) < \infty : \\
& \quad \text{(Green)}, \quad \text{fi} \quad ( \quad ) \\
& [\mathcal{T}; o] \\
& (1) \quad \mathcal{O}(J) = \langle \mathcal{O}; \mathbb{1}_J(\mathcal{A}_{\mathcal{T}}) \mathcal{O} \rangle \quad J \subseteq \mathbb{R} : \\
& \quad \text{(Green)} \quad \sup_{\lambda \in I, \eta_0 > 0} \mathbb{E}(|\mathcal{G}^\gamma(o; o)|^2) < \infty. \\
& \quad \mathbb{P} \dots [\mathcal{T}; o], \quad \frac{1}{\pi} \text{Im } \mathcal{G}^{\lambda + i0}(o; o), \quad \text{(Green)} \quad \mathbb{P} \dots \\
& \quad \mathcal{A}_{\mathcal{T}} \quad I, \quad 12 \quad \mathcal{I}. \\
& \quad \text{(Green)}, \quad s < 0, \quad \mathcal{I} \\
& \quad \mathcal{I} \quad \mathcal{I}_1 \subset \mathcal{I}. \quad G_N \quad \Gamma_N \setminus \widetilde{G_N} \quad \widetilde{G_N} \\
& \quad (\quad) \quad \mathcal{G}_N. \quad \tilde{x}; \tilde{y} \quad \widetilde{G_N}, \quad \in \mathbb{C} \setminus \mathbb{R}, \quad \mathcal{A}_N \quad \widetilde{G_N} \\
& (1) \quad \tilde{g}_N^\gamma(\tilde{x}; \tilde{y}) = \langle \tilde{x}; (\tilde{\mathcal{A}}_N - \cdot)^{-1} \tilde{y} \rangle_{\ell^2(\widetilde{G_N})} :
\end{aligned}$$

**Theorem 12** ( ). Assume that  $(G_N; W_N)$  satisfies **(BSCT)**, **(EXP)** and **(Green)**.

Call  $(\lambda_j^{(N)})_{j=1}^N$  the eigenvalues of  $\mathcal{A}_N$  on  $\mathcal{V}^2(V_N)$ , and let  $(\gamma_j^{(N)})_{j=1}^N$  be a corresponding orthonormal eigenbasis.

For each  $N$ , let  $a = a_N$  be a function on  $V_N$  with  $\sup_N \sup_{x \in V_N} |a_N(x)| \leq 1$ .

Then

$$\lim_{\eta_0 \downarrow 0} \lim_{N \rightarrow +\infty} \frac{1}{N} \sum_{\lambda_j^{(N)} \in I_1} \left| \sum_{x \in V_N} a(x) |\gamma_j^{(N)}(x)|^2 - \sum_{x \in V_N} a(x) \frac{N}{\lambda_j^{(N)} + i\eta_0} \gamma_j^{(N)}(x) \right| = 0;$$

for some family of probability measures  $\gamma_N$  on  $V_N$ , indexed by a parameter  $\gamma \in \mathbb{C} \setminus \mathbb{R}$ , defined as follows :

$$\gamma_N(x) = \frac{\text{Im } \tilde{g}_N^\gamma(\tilde{x}; \tilde{x})}{\sum_{y \in V_N} \text{Im } \tilde{g}_N^\gamma(\tilde{y}; \tilde{y})}.$$

$$(1) \quad \frac{1}{N} \left| \left\{ j^{(N)} \in I_1 : \left| \sum_{x \in V_N} a(x) |j^{(N)}(x)|^2 - \sum_{x \in V_N} a(x) \frac{N}{\lambda_j^{(N)} + i\eta_0} (x) \right| > \right\} \right| \xrightarrow{N \rightarrow +\infty, \eta_0 \downarrow 0} 0 :$$

**Corollary 13.** *For any  $\alpha \in (0;1)$ , there exists  $c_\alpha > 0$  such that for any  $\Lambda_N \subset V_N$  with  $|\Lambda_N| \geq N$ , we have*

$$(20) \quad \frac{1}{N} \sum_{x \in V_N} F \left( N^{\frac{N}{\lambda + i\eta_0}}(x) \right) \xrightarrow{N \rightarrow +\infty} \mathbb{E} \left( F \left( \frac{\text{Im } \mathcal{G}^{\lambda + i\eta_0}(o; o)}{\mathbb{E}(\text{Im } \mathcal{G}^{\lambda + i\eta_0}(o; o))} \right) \right) :$$

**Remark 14.** *The results proven in [1] actually hold for more general Schrödinger operators than adjacency matrices : one can consider weighted Laplacians (with conductances on the edges) and add a potential; in other words, on each  $G_N$ , we can consider a discrete Schrödinger operator  $\mathcal{H}_N$ . The limiting object in assumption **(BSCT)** is now a random*

rooted tree  $[\mathcal{T}; o]$  endowed with a random Schrödinger operator  $\mathcal{H}$ . Assumption **(Green)** has to be modified, replacing the adjacency matrix  $\mathcal{A}$  by the operator  $\mathcal{H}$ . Similarly, in the statement of the theorem, the Green functions  $\tilde{g}_N$  to be considered are those of  $\mathcal{H}_N$  lifted to the universal cover  $\widetilde{G}_N$ .

**Remark 15.** In particular, our result applies to the case where the limiting system  $([\mathcal{T}; o]; \mathcal{H})$  is  $\mathcal{T} = \mathfrak{X}$  (the  $(q+1)$ -regular tree) with an arbitrary origin  $o$ , and  $\mathcal{H} = \mathcal{H}_\epsilon = \mathcal{A} + \mathcal{W}$  where  $\mathcal{W}$  is a random real-valued potential on  $\mathfrak{X}$ . More precisely the values  $\mathcal{W}(x)$  ( $x \in \mathfrak{X}$ ) are i.i.d. random variables of common law  $\mu$ . This is known as the Anderson model on  $\mathfrak{X}$ . It was shown by A. Klein [Kl80] that the spectrum of  $\mathcal{H}_\epsilon$  is a.s. purely absolutely continuous on  $I = (-2\sqrt{q} + \epsilon; 2\sqrt{q} - \epsilon)$ , provided  $\epsilon$  is small enough (depending on  $\mu$ ). This just assumes a second moment on  $\mu$ . Under stronger regularity assumptions on  $\mu$ , one can show that Assumption **(Green)** holds on  $I$  (see [11], following Aizenman-Warzel [2]). Examples of sequences of expander regular graphs  $G_N$  with discrete Schrödinger operators  $\mathcal{H}_N$  converging to  $([\mathfrak{X}; o]; \mathcal{H}_\epsilon)$  are given in [10].

**Remark 16.** Examples of sequences of  $(G_N)$  satisfying our three assumptions were investigated in [11]. In the examples considered there, the limiting trees  $\mathcal{T}$  are  $\mathfrak{X}$  or  $\mathfrak{Y}$ ; roughly speaking, those are trees where the local geometry can only take a finite number of values. If  $\mathcal{A}$  is the adjacency matrix of such a tree, we showed in [11] that the spectrum of  $\mathcal{A}$  is a finite union of closed intervals, and that there are a finite number of points  $y_1; \dots; y_\ell$  in  $\mathbb{R}$  such that Assumption **(Green)** holds on any  $I$  of the form  $\mathbb{R} \setminus ([y_1 - \epsilon; y_1 + \epsilon] \cup \dots \cup [y_\ell - \epsilon; y_\ell + \epsilon])$  (for any  $\epsilon > 0$ ). We showed – extending Remark 15 – that on such trees, Assumption **(Green)** remains true after adding a small random potential to  $\mathcal{A}_\mathcal{T}$ . Finally, we showed the existence of sequences  $(G_N)$  converging to  $\mathcal{T}$  and satisfying the **(EXP)** condition.

PEPEC AND LINK IHO HE OK

1. Random regular graphs. [11]

deterministic  $(G_N)$  a.

random random

fi, 1, N

(1, 13, 1)

adjacency matrix  $a_N : \{1; \dots; N\} \rightarrow \mathbb{R}$



$$\mathcal{N}(0; \frac{2}{j}) \quad 0 \leq j \leq 1. \quad \sqrt{N} \quad j(x) \quad x \quad \{1; \dots; N\} \quad R \geq 0, \quad (\sqrt{N} \quad j(y))_{y, d(y, x) \leq R} \quad G_N, \quad B_{\mathfrak{X}}(0; R).$$

## 2. From graph Laplacians to Hecke operators.

11. From graph Laplacians to Hodge operators

Let  $(g_1, \dots, g_k) \in \Delta_{\mathbb{S}^2}$  be a free subgroup of  $SO(3)$ .

$$T_k f(x) = \sum_{j=1}^k (f(g_j x) + f(g_j^{-1} x))$$

$$\Delta_{\mathbb{S}^2}.$$

**Theorem 19** ( ). Assume that  $g_1, \dots, g_k$  generate a free subgroup of  $SO(3)$ .

For each  $\ell$ , let  $(\psi_j^{(\ell)})_{j=1}^{2\ell+1}$  be an orthonormal family of eigenfunctions of  $-\Delta_{\mathbb{S}^2}$  of eigenvalue  $\ell(\ell+1)$ , that are also eigenfunctions of  $T_k$ .

Then for any continuous function  $a$  on  $\mathbb{S}^2$ , we have

$$\frac{1}{2\ell+1} \sum_{j=1}^{2\ell+1} \left| \int_M a(x) |\psi_j^{(\ell)}(x)|^2 d\text{Vol}(x) - \int_M a(x) d\text{Vol}(x) \right|^2 \xrightarrow{\ell \rightarrow \infty} 0.$$

11.  $T_k$  is not a pseudodifferential operator, so the argument sketched above to show that the basis  $(Y_\ell^m)_{\ell \geq 0, |m| \leq \ell}$  could not satisfy quantum ergodicity does not apply here.

**Remark 20.**  $T_k$  is not a pseudodifferential operator, so the argument sketched above to show that the basis  $(Y_\ell^m)_{\ell \geq 0, |m| \leq \ell}$  could not satisfy quantum ergodicity does not apply here.

**Remark 21.** We note that for very special choices of rotations – rotations that correspond to norm  $n$  elements in an order in a quaternion division algebra, the operators  $T_k$  are called Hecke operators. It has been conjectured by Böcherer, Sarnak, and Schulze-Pillot [2] that such joint eigenfunctions satisfy the much stronger quantum unique ergodicity property. This conjecture is still open.

3. [1] 2000,  $(n)$  infinite  $(\cdot, \cdot)$  almost-every ergodic component,  $M$ , 10 one



11

### 3. Quantum ergodicity on Riemann surfaces of high genus.

11 12

11

**Theorem 22** ( $M \rightarrow \infty$ ). Let  $(S_N)$  be a sequence of hyperbolic surfaces, whose genus (equivalently, volume) goes to  $\infty$ .

**(EXP)** Assume the first eigenvalue  $\lambda_1(N)$  of  $-\Delta$  on  $S_N$  is bounded away from 0 as  $N \rightarrow \infty$ .

**(BSH)** Assume there are few short geodesics; in other words,  $(S_N)$  converges in the Benjamini-Schramm sense to the hyperbolic disc: for any  $R > 0$ ,

$$\lim_{N \rightarrow +\infty} \frac{\text{Vol}\{x \in S_N : r(x) < R\}}{\text{Vol}(S_N)} = 0$$

where  $r(x)$  means the injectivity radius at  $x$ .

Fix an interval  $I \subset (0; +\infty)$ .

Let  $(\phi_i^{(N)})$  be an orthonormal basis of eigenfunctions of the Laplacian on  $S_N$ .

Let  $a = a_N : S_N \rightarrow \mathbb{C}$  be such that  $|a(x)| \leq 1$  for all  $x \in S_N$ . Then

$$\lim_{N \rightarrow +\infty} \frac{1}{\text{Vol}(S_N)} \sum_{\lambda_i(N) \in I} \left| \int_{S_N} a(x) |\phi_i^{(N)}(x)|^2 dx - \langle a \rangle \right|^2 = 0$$

where  $\langle a \rangle = \frac{1}{\text{Vol}(S_N)} \int_{S_N} a(x) dx$ .

(1; +\infty) L^2- S\_N

22

random

random

1

(\phi\_i^{(N)})

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