



In quantum mechanics, the position  $q$  and momentum  $p$  operators satisfy the canonical commutation relation  $qp - pq = i\hbar I$ , where  $\hbar$  is the reduced Planck constant. The Hamiltonian  $H$  is given by  $H = \frac{p^2}{2m} + V(q)$ , where  $m$  is the mass and  $V$  is the potential energy. The time evolution of an operator  $A$  is governed by the Heisenberg equation of motion  $i\hbar \frac{dA}{dt} = [A, H]$ . The energy eigenvalues  $E_n$  are determined by the eigenvalue equation  $H\psi_n = E_n\psi_n$ . The wave function  $\psi(x, t)$  satisfies the Schrödinger equation  $i\hbar \frac{\partial}{\partial t} \psi = \left( -\frac{\hbar^2}{2m} \Delta + V \right) \psi$ .

The wave function  $\psi(x, t)$  is a function on  $\mathbb{R}^3$  and is normalized such that  $\int_{\mathbb{R}^3} |\psi(x, t)|^2 dx = 1$ . The energy eigenvalues  $E_n$  are real numbers. The wave function  $\psi(x, t)$  can be expanded in terms of energy eigenfunctions  $\psi_n(x)$  as  $\psi(x, t) = \sum_n c_n \psi_n(x) e^{-itE_n/\hbar}$ .

$$\psi(x, t) = \sum_n c_n \psi_n(x) e^{-itE_n/\hbar}$$

The coefficients  $c_n$  are determined by the initial condition  $\psi(x, 0) = \psi_0(x)$ . The energy eigenfunctions  $\psi_n(x)$  are orthogonal and form a complete basis for the Hilbert space  $L^2(\mathbb{R}^3)$ . The time evolution of the wave function is unitary, meaning that the norm of the wave function is preserved. The probability density  $|\psi(x, t)|^2$  is conserved in time, as shown by the continuity equation  $\frac{\partial}{\partial t} |\psi|^2 + \nabla \cdot \mathbf{j} = 0$ , where  $\mathbf{j}$  is the probability current.

$H(x; t) = H = E, E \in \mathbb{R}$ ,  
 (1).

$\mathbb{R}^2$ ,

*chaotic*,

$2$ .

103,

*arithmetic*

$\mathbb{R}$

$2$ ,

$3$ ,

$4$ ,

$3.4$ ,

$$\begin{aligned}
 & \int_M |\Delta_k \psi| \, \text{Vol} \sim \int_M |\psi| \, \text{Vol} \\
 & \int_M |\Delta_k \psi|^2 \, \text{Vol} \sim \int_M |\psi|^2 \, \text{Vol} \\
 & \int_M |\Delta_k \psi| \, \text{Vol} \sim \int_M |\psi| \, \text{Vol}
 \end{aligned} \tag{4.1}$$

## 2. HIGH FREQUENCY ENCODING AND DELOCALIZATION

$$\begin{aligned}
 & \int_M |\Delta_k \psi| \, \text{Vol} \sim \int_M |\psi| \, \text{Vol} \\
 & \int_M |\Delta_k \psi|^2 \, \text{Vol} \sim \int_M |\psi|^2 \, \text{Vol} \\
 & \int_M |\Delta_k \psi| \, \text{Vol} \sim \int_M |\psi| \, \text{Vol}
 \end{aligned}$$

$(M; g)$  is a Riemannian manifold of dimension  $d$ . We consider the Laplacian  $\Delta_k = -\hbar^2 \Delta$  on  $L^2(M; \text{Vol})$ . For  $k \in \mathbb{N}$ , we consider the eigenvalues  $\lambda_k$  and eigenfunctions  $\psi_k$  of  $\Delta_k$ . The eigenvalues are given by  $\lambda_k = -k^2$ . The eigenfunctions  $\psi_k$  are localized near the boundary of  $M$ . As  $k \rightarrow +\infty$ , the eigenfunctions  $\psi_k$  become more and more localized near the boundary of  $M$ . This is the semiclassical limit.

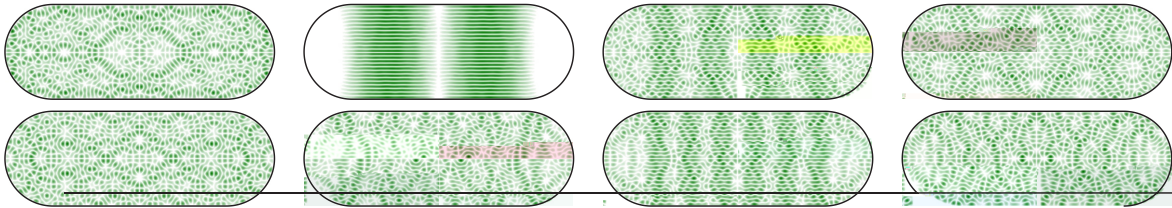


FIGURE 1.  $|\psi_n(x; y)|^2$  for  $n = 319$ .

### 2.1. The role of the geodesic flow.

$$\begin{aligned}
 & -\hbar^2 \Delta \psi = E \psi \\
 & \int_M |\Delta_k \psi|^2 \, \text{Vol} \sim \int_M |\psi|^2 \, \text{Vol}
 \end{aligned}$$



(2),  $\| \lambda_k \|_\infty \sim k^{-1/2}$ .  $\int_{\mathbb{R}^d} e^{i\langle x, \lambda \rangle} \lambda \, d\mu(\lambda) = 1$ .  $\| \lambda \|_{L^2} = 1$ .  $L^\infty$ .

2.2.  $L^p$ -norms as measures of delocalization ?

$\lambda = \lambda$ ,  $\| \lambda \|_{L^2} = 1$ .  $L^\infty$ .

**Theorem 1.** (known as Hörmander’s bound)

$$\| \lambda_k \|_\infty = O(k^{-(d-1)/4});$$

$\| \lambda \|_{L^p} \sim k^{-\mu(p)}$ ,  $2 \leq p \leq +\infty$ ,

**Theorem 2** ( [10], 10 ).

$$\| \lambda \|_{L^p} = O(k^{-\mu(p)})$$

where

- $\mu(p) = d\left(\frac{1}{2} - \frac{1}{p}\right) - \frac{1}{2}$  for  $\frac{2(d+1)}{(d-1)} \leq p \leq +\infty$ ;
- $\mu(p) = \frac{d-1}{2}\left(\frac{1}{2} - \frac{1}{p}\right)$  for  $2 \leq p \leq \frac{2(d+1)}{(d-1)}$ .

$M$ .  $\mathbb{R}^d$ .  $p_c = \frac{2(d+1)}{(d-1)}$ .  $p \geq p_c$  ( $L^p$ -norms),  $p \leq p_c$  ( $L^p$ -norms),  $p > p_c$ ,  $L^p$ -norms.

$$\mathcal{L}_x = \{v \in S_x^*M; \exists t > 0; g^t(x; v) \in S_x^*M\}:$$

**Theorem 3** ( [10], 10, 110 ).

- Assume there exists a subsequence  $n_k \rightarrow +\infty$  and  $C > 0$  such that  $\| \lambda_{n_k} \|_\infty \geq C k^{-\frac{\mu(\infty)}{2}}$ . Then there exists  $x$  such that  $\mathcal{L}_x > 0$ ;
- If  $M$  is real analytic, the existence of such subsequence  $\lambda_{n_k}$  is sufficient to the existence of  $x$  such that  $\mathcal{L}_x = S_x^*M$ , and the first return map  $\tau_x : S_x^*M \rightarrow S_x^*M$  possesses an absolutely continuous invariant probability measure. (Moreover, in

that case, there exists  $t_0 > 0$  such that  $g^{t_0}(x; v) \in S_x^*M$  for all  $v \in S_x^*M$ , that is, there is a common return time).

- If  $M$  is real analytic and  $\dim M = 2$ , the existence of such subsequence  $\lambda_{n_k}$  is due to the existence of  $x \in M$  and  $t_0 > 0$  such that  $g^{t_0}(x; v) = (x; v)$  for all  $v \in S_x^*M$ .

Anosov property,

$$\|\lambda\|_{L^\infty} = o\left(\frac{\mu(\infty)}{2}\right) \quad (O(1) \text{ term}).$$

**Theorem 4.** (i) If  $d = 2$  and  $M$  has no conjugate points, or if  $d \geq 2$  and  $M$  has non-positive sectional curvature, for  $p = +\infty$ ,

$$\|\lambda\|_{L^p} = O\left(\frac{\mu(p)}{\sqrt{\log}}\right):$$

(i') Statement (i) actually holds if  $M$  has no conjugate points, for all  $d \geq 2$ .

(ii) (i) holds for all  $p > p_c$ .

(iii) If  $M$  has non-positive sectional curvature, for  $p < p_c$ , there exists  $(p; d) > 0$  such that

$$\|\lambda\|_{L^p} = O\left(\frac{\mu(p)}{(\log)^{\sigma(p,d)}}\right)$$

(iv) Statement (iii) still holds for  $p = p_c$ .

### 2.3. The Shnirelman theorem and the Quantum Unique ergodicity conjecture .

Let  $(M, g)$  be a compact Riemannian manifold with volume  $Vol(M)$ . For a function  $k$  on  $M$ , define the  $L^2$  norm  $\|k\|_{L^2} = \left(\int_M |k(x)|^2 dVol(x)\right)^{1/2}$ . The Shnirelman theorem states that for any function  $k$  on  $M$ ,  $\|k\|_{L^2} \rightarrow 0$  as  $k \rightarrow +\infty$ .

Let  $(x_0; 0) \in S^*M$ ,  $L^1$ -continuous function  $a : S^*M \rightarrow \mathbb{R}$ ,  $T \rightarrow +\infty$

$$\left( \int_{S^*M} a dL - \frac{1}{T} \int_0^T \int_{S^*M} a \circ g^t(x_0; 0) dt \right) \xrightarrow{T \rightarrow +\infty} 0$$

for every  $(x_0; 0) \in S^*M$ .

### Quantum Ergodicity Theorem (Shnirelman theorem).

**Theorem 5** ([10, 40, 11]). Let  $(M; g)$  be a compact Riemannian manifold, with the metric normalized so that  $\text{Vol}(M) = 1$ . Call  $\Delta$  the Laplace-Beltrami operator on  $M$ . Assume that the geodesic flow of  $M$  is ergodic with respect to the Liouville measure. Let  $(\phi_k)_{k \in \mathbb{N}}$  be an orthonormal basis of  $L^2(M; g)$  made of eigenfunctions of the Laplacian

$$\Delta \phi_k = -\lambda_k \phi_k; \quad \lambda_k \xrightarrow{k \rightarrow \infty} +\infty.$$

Let  $a$  be a continuous function on  $M$ . Then

$$(3) \quad \frac{1}{N(\lambda)} \sum_{k, \lambda_k \leq \lambda} \left| \langle \phi_k; a \phi_k \rangle_{L^2(M)} - \int_M a(x) d\text{Vol}(x) \right| \xrightarrow{\lambda \rightarrow +\infty} 0$$

where the normalizing factor is  $N(\lambda) = |\{k; \lambda_k \leq \lambda\}|$ :

$$\langle \phi_k; a \phi_k \rangle_{L^2(M)} = \int_M a(x) |\phi_k(x)|^2 d\text{Vol}(x).$$

**Remark 6.** The Cesaro limit (3) implies that there exists a subset  $S \subset \mathbb{N}$  of density 1 such that

$$(4) \quad \langle \phi_k; a \phi_k \rangle \xrightarrow{n \rightarrow +\infty, n \in S} \int_M a(x) d\text{Vol}(x):$$

In addition, using the fact that the space of continuous functions is separable, one can actually find  $S \subset \mathbb{N}$  of density 1 such that (4) holds for every  $a \in C^0(M)$ . In other words, the sequence of measures  $(|\phi_k(x)|^2 d\text{Vol}(x))_{n \in S}$  converges weakly to the uniform measure  $d\text{Vol}(x)$ .

Actually, the full statement of the theorem says that there exists a subset  $S \subset \mathbb{N}$  of density 1 such that

$$(5) \quad \langle \phi_k; A \phi_k \rangle \xrightarrow{n \rightarrow +\infty, n \in S} \int_{S^*M} \sigma^0(A) dL$$

for every pseudodifferential operator  $A$  of order 0 on  $M$ . On the right-hand side,  $\sigma^0(A)$  is the principal symbol of  $A$ , that is a function on the unit cotangent bundle  $S^*M$ . Equation 4 corresponds to the case where  $A$  is the operator of multiplication by the function  $a$ .





2.1.  $(n_k)$ .  $g_{\#}^t = \dots$   $t \in \mathbb{R}$ .  $\rightarrow +\infty$

**Theorem 7.3** Assume  $M$  is a compact Riemannian manifold with negative sectional curvature. Assume  $\langle n_k; A_{n_k} \rangle$  converges to  $\int_{S^*M} \chi(A) d\mu$  for all  $A$ . Then  $\mu$  has positive entropy.

$h_{KS}(\mu) = \int \sum_{j=1}^r \chi_j(\mu) d\mu$   
 $h_{KS}(\mu) = 0$   
 $\mu \mapsto h_{KS}(\mu) = \int \chi(\mu) d\mu$   
 $h_{KS}(\mu_1 + (1-\alpha)\mu_2) = \alpha h_{KS}(\mu_1) + (1-\alpha)h_{KS}(\mu_2)$ ,  $\alpha \in [0,1]$   
 $(\mu_1 - \mu_2) \in \mathcal{R}$   
 $(\mu) \quad h_{KS}(\mu) \leq \int \left( \sum_{j=1}^r \chi_j(\mu) \right) d\mu$   
 $(\mu) \quad h_{KS}(\mu) \leq d-1$   
 $L$

**Corollary 1.** Let  $\Gamma \subset S^*M$  be the union of all points lying on a periodic trajectory on the geodesic flow (recall, if  $M$  has negative curvature, there are countably many periodic geodesics). Let  $\mu$  be as in Theorem 7. Then  $\mu(\Gamma) < 1$ .

$\mu$  strongly scarred

Let  $\Gamma$  be a partial scar,  $\Gamma \subset S^*M$ ,  $(\Gamma) > 0$ .

**Corollary 2.** *The support of  $\mu_\Gamma$  has Hausdorff dimension  $> 1$ .*

$\dim_{\text{Haus}}(\text{supp } \mu_\Gamma) \geq 1$

**Theorem 8** ( ). Assume  $M$  is a compact Riemannian manifold of dimension  $d$ , with constant sectional curvature  $-1$ . Then  $\mu_\Gamma$  has entropy greater than  $\frac{d-1}{2}$ .

$h_{KS}(\mu_\Gamma) > \frac{d-1}{2}$

**Corollary 3.** *The support of  $\mu_\Gamma$  has Hausdorff dimension  $\geq d$ .*

$\dim_{\text{Haus}}(\text{supp } \mu_\Gamma) \geq d$

**Corollary 4.** Let  $\Gamma \subset S^*M$  be the union of all points lying on a closed trajectory on the geodesic flow. Let  $\mu_\Gamma$  be as in Theorem 7, with  $M$  of constant negative curvature. Then  $(\Gamma) \leq 1=2$ .

$h_{KS}(\mu_\Gamma) = 0$ ,  $h_{KS}(\mu_\Gamma) = (1 - \frac{1}{2}) h_{KS}(\mu_\Gamma)$ ,  $\Gamma$ ,  $\geq \frac{d-1}{2}$ ,  $\leq 1=2$ .

$\frac{1}{2} \int_{S^*M} \sum_{j=1}^{d-1} j^2 d\mu_\Gamma$ ,  $d=2$ ,  $1=2$

**Theorem 9** ( ).  $\mu_\Gamma$  has full support, that is,  $(\Omega) > 0$  for any non-empty open set  $\Omega \subset S^*M$ .

Let  $\Omega$  be a domain in  $\mathbb{R}^n$  with boundary  $Z$ . Consider the Dirichlet problem for the Laplacian on  $\Omega$ . The eigenvalues  $\lambda_j$  of the Laplacian on  $\Omega$  are real and non-negative. The eigenfunctions  $\phi_j$  form an orthonormal basis for  $L^2(\Omega)$ . The eigenvalues  $\lambda_j$  are asymptotically distributed according to Weyl's law.

2. . Some questions on non-compact manifolds.

Let  $M$  be a non-compact manifold. Consider the Laplacian on  $M$ . The spectrum of the Laplacian on  $M$  is continuous. The eigenfunctions  $\phi_j$  are not square-integrable. The eigenvalues  $\lambda_j$  are asymptotically distributed according to Weyl's law.

2. .1. Absolutely continuous spectrum .

Let  $M$  be a non-compact manifold. Consider the Laplacian on  $M$ . The spectrum of the Laplacian on  $M$  is absolutely continuous. The eigenfunctions  $\phi_j$  are not square-integrable. The eigenvalues  $\lambda_j$  are asymptotically distributed according to Weyl's law.

2. .2. Large frequency delocalization on non-compact manifolds.

Let  $M$  be a non-compact manifold. Consider the Laplacian on  $M$ . The eigenfunctions  $\phi_j$  are delocalized as  $\lambda_j \rightarrow +\infty$ .



3.1. Overview of the problem.  $G = (V; E)$ . *regular*

3.1. Overview of the problem.

localized, delocalized  $G = (V; E)$ .  
 $\Leftrightarrow$   $(\cdot, \cdot)$ ,  $I \subset \mathbb{R}$ ,

- *spectral localization* :
- *exponential localization* :
- *dynamical localization* :
- *spectral delocalization* :
- *ballistic transport* :
- *spatial delocalization*.

$(G_N)$   $N \rightarrow \infty$ .  $(G_N)$   $N \times N$

$\sum_{x=1}^N |j(x)|^2$   $1=N$

- $\infty$  norms :  $\|j\|_\infty$
- $p$  norms :  $\|j\|_p$   $N^{1/p-1/2}$   $f(p) \neq 1=p-1=2$ .
- Scarring :  $(\sum_{x \in \Lambda} |j(x)|^2 \geq 1 - \epsilon)$
- Quantum ergodicity :  $a : \{1; \dots; N\} \rightarrow \mathbb{C}$

most  $j$ .  $\forall$  all  $j$ , quantum unique ergodicity.  
 almost sure random random  
 2, 3, 33, 1, 1, 13, 14 4.1. ff

3.2. Entropy.

$(q + 1)$ -  
 $(q + 1)$ -  
 $(G_N)_{N \in \mathbb{N}} = (V_N; E_N)$ .

( ) 
$$\mathcal{A}_N f(x) = \sum_{x \sim y} f(y)$$

( ) 
$$\Delta_N f(x) = \sum_{x \sim y} (f(y) - f(x)) :$$

(10) 
$$\mathcal{A}_N - (q + 1)I = \Delta_N :$$

**Theorem 10** ( ). Let  $(G_N)$  be a sequence of  $(q + 1)$ -regular graphs (with  $q$  fixed),  $G_N = (V_N; E_N)$  with  $V_N = \{1; \dots; N\}$ . Assume that<sup>1</sup> there exists  $c > 0$ ;  $\delta > 0$  such that, for any  $k \leq c \ln N$ , for any pair of vertices  $x; y \in V_N$ ,

(11) 
$$|\{\text{paths of length } k \text{ in } G_N \text{ from } x \text{ to } y\}| \leq q^{k(\frac{1-\delta}{2})} :$$

Fix  $\delta > 0$ . Then, if  $f$  is an eigenfunction of the discrete Laplacian on  $G_N$  and if  $\Lambda \subset V_N$  is a set such that

$$\sum_{x \in \Lambda} |f(x)|^2 \geq \sum_{x \in V_N} |f(x)|^2 ;$$

then  $|\Lambda| \geq N^\alpha$  — where  $\alpha > 0$  is given as an explicit function of  $\delta$ ; and  $c$ .

$$H_N(f) = -\frac{1}{\log N} \sum_x |f(x)|^2 \ln |f(x)|^2$$

1

$\geq c \ln N$





**Theorem 11** ( ). Let  $(G_N) = (V_N; E_N)$  be a sequence of  $(q + 1)$ -regular graphs with  $|V_N| = N$ . Assume that  $(G_N)$  satisfies **(BST)** and **(EXP)**.

Let  $(\varphi_1^{(N)}; \dots; \varphi_N^{(N)})$  be an orthonormal basis of eigenfunctions of  $\mathcal{A}_N$  in  $\ell^2(V_N)$ .

Let  $a_N : V_N \rightarrow \mathbb{C}$  be a sequence of functions such that  $\sup_N \sup_{x \in V_N} |a_N(x)| \leq 1$ : Define  $\langle a_N \rangle = \frac{1}{N} \sum_{x \in V_N} a_N(x)$ .

Then

$$\begin{aligned}
 & \frac{1}{N} \sum_{j=1}^N \left| \langle \varphi_j^{(N)}; a_N \varphi_j^{(N)} \rangle_{\ell^2(V_N)} - \langle a_N \rangle \right|^2 \xrightarrow{N \rightarrow +\infty} 0: \\
 & \text{ , } > 0, \\
 (13) \quad & \frac{1}{N} \left| \left\{ j \in [1; N]; \left| \langle \varphi_j^{(N)}; a_N \varphi_j^{(N)} \rangle_{\ell^2(V_N)} - \langle a_N \rangle \right| > \right\} \right| \xrightarrow{N \rightarrow +\infty} 0: \\
 & \left\langle \varphi_j^{(N)}; a_N \varphi_j^{(N)} \right\rangle_{\ell^2(V_N)} = \frac{1}{N} \sum_{x \in V_N} a_N(x) |\varphi_j^{(N)}(x)|^2, \\
 & \sum_{x \in V_N} | \varphi_j^{(N)}(x) |^2 = \frac{1}{N} \sum_{x \in V_N} 1 = 1, \\
 & \left( \sum_{x \in V_N} | \varphi_j^{(N)}(x) |^2 \right) = \frac{1}{N} \sum_{x \in V_N} 1 = 1. \tag{13}
 \end{aligned}$$

### 3.4. Non-regular graphs : from spectral to spatial delocalization.

regular quantum ergodicity theorem spectral delocalization spatial delocalization

“If a large finite system is close (in the Benjamini-Schramm topology) to an infinite system having purely absolutely continuous spectrum in an interval  $I$ , then the eigenfunctions (with eigenvalues lying in  $I$ ) of the finite system satisfy quantum ergodicity.”

$$\begin{aligned}
 (G_N)_{N \in \mathbb{N}} \quad & V_N \xrightarrow{N \rightarrow \infty} \mathbb{Z}^d, \quad \mathcal{A}_N : \mathbb{C}^{V_N} \rightarrow \mathbb{C}^{V_N} \\
 & \mathcal{A}_N f(v) = \sum_{w \sim v} f(w); \\
 & V \sim W \quad \mathcal{A}_N \quad \left( \frac{\mathbb{Z}}{N} / \frac{\mathbb{Z}}{N} \right) \quad N \rightarrow +\infty.
 \end{aligned}$$



$$\mathcal{T} \begin{matrix} v, w \in \mathcal{T} \\ v \sim w, \\ \mathcal{T}^{(v|w)} \\ \mathcal{G}^{(v|w)}(\cdot; \cdot; \cdot) \\ \gamma_w(v) := -\mathcal{G}^{(v|w)}(v; v; \cdot). \end{matrix} \begin{matrix} \mathcal{A}_{\mathcal{T}} \end{matrix}$$

(Green)  $I, s > 0$

$$\sup_{\lambda \in I, \eta_0 \in (0,1)} \mathbb{E} \left( \sum_{y: y \sim o} |\text{Im } o^{\lambda+i\eta_0}(y)|^{-s} \right) < \infty :$$

(Green), fi

$[\mathcal{T}; o]$

$$(1) \quad o(J) = \langle o; \mathbb{1}_J(\mathcal{A}_{\mathcal{T}}) o \rangle \quad J \subseteq \mathbb{R} :$$

$$(Green) \quad \sup_{\lambda \in I, \eta_0 > 0} \mathbb{E}(|\mathcal{G}^\gamma(o; o)|^2) < \infty.$$

$$I, \frac{1}{\pi} \text{Im } \mathcal{G}^{\lambda+i0}(o; o), (Green) \mathbb{P}^{\cdot \cdot}$$

$$A_{\mathcal{T}}, I, s < 0,$$

$$I, (Green), \text{fi } I_1$$

$$I_1 \subset I, \tilde{x}, \tilde{y} \in \mathbb{C} \setminus \mathbb{R}, \tilde{A}_N, \tilde{G}_N$$

$$(1) \quad \tilde{g}_N^\gamma(\tilde{x}; \tilde{y}) = \langle \tilde{x}; (\tilde{A}_N - \cdot)^{-1} \tilde{y} \rangle_{\ell^2(\tilde{G}_N)} :$$

**Theorem 12** ( ). Assume that  $(G_N; W_N)$  satisfies (BSCT), (EXP) and (Green).

Call  $(\lambda_j^{(N)})_{j=1}^N$  the eigenvalues of  $\mathcal{A}_N$  on  $\ell^2(V_N)$ , and let  $(\gamma_j^{(N)})_{j=1}^N$  be a corresponding orthonormal eigenbasis.

For each  $N$ , let  $a = a_N$  be a function on  $V_N$  with  $\sup_N \sup_{x \in V_N} |a_N(x)| \leq 1$ .

Then

$$\lim_{\eta_0 \downarrow 0} \lim_{N \rightarrow +\infty} \frac{1}{N} \sum_{\lambda_j^{(N)} \in I_1} \left| \sum_{x \in V_N} a(x) |\lambda_j^{(N)}(x)|^2 - \sum_{x \in V_N} a(x) \lambda_j^{(N)+i\eta_0}(x) \right| = 0 ;$$

for some family of probability measures  $\gamma_N$  on  $V_N$ , indexed by a parameter  $\gamma \in \mathbb{C} \setminus \mathbb{R}$ , defined as follows :

$$\gamma_N(x) = \frac{\text{Im } \tilde{g}_N^\gamma(\tilde{x}; \tilde{x})}{\sum_{y \in V_N} \text{Im } \tilde{g}_N^\gamma(\tilde{y}; \tilde{y})} :$$

Here,  $\tilde{x} \in \widetilde{G}_N$  is a lift of  $x \in V_N$ .

$$(1) \quad \frac{1}{N} \left\{ \binom{N}{j} \in I_1 : \left| \sum_{x \in V_N} a(x) | \binom{N}{j}(x) |^2 - \sum_{x \in V_N} a(x) \binom{N}{\lambda_j^{(N)} + i\eta_0}(x) \right| > \right\} \xrightarrow{N \rightarrow +\infty, \eta_0 \downarrow 0} 0 :$$

$a_N = \mathbb{1}_{\Lambda_N}, \quad \Lambda_N \subset V_N, \quad |\Lambda_N| \approx N, \quad \in (0;1).$

$\mathcal{G}^\gamma(\tilde{x}; \tilde{y}) = \mathcal{G}^\gamma(o; o), \quad \tilde{x} \in \widetilde{G}_N, \quad \mathcal{G}^\gamma(\tilde{x}; \tilde{y}) = \mathcal{G}^\gamma(o; o) + \mathcal{O}(\frac{1}{N})$

$\|\mathbb{1}_{\Lambda_N} \binom{N}{j}\|^2 \approx \binom{N}{j}, \quad \Lambda_N \subset V_N$

$c_\alpha > 0, \quad \Lambda_N \subset V_N, \quad |\Lambda_N| \geq N,$

$$(1) \quad \inf_{\eta_0 \in (0,1)} \liminf_{N \rightarrow \infty} \inf_{\lambda \in I} \binom{N}{\lambda + i\eta_0}(\Lambda_N) \geq 2c_\alpha :$$

(1),

**Corollary 13.** For any  $\in (0;1)$ , there exists  $c_\alpha > 0$  such that for any  $\Lambda_N \subset V_N$  with  $|\Lambda_N| \geq N$ , we have

$$\frac{1}{N} \# \left\{ \binom{N}{j} \in I : \|\mathbb{1}_{\Lambda_N} \binom{N}{j}\|^2 < c_\alpha \right\} \xrightarrow{N \rightarrow +\infty} 0 :$$

$\|\mathbb{1}_{\Lambda_N} \binom{N}{j}\|^2 \approx \binom{N}{j}, \quad \Lambda_N \subset V_N$

$\|\mathbb{1}_{\Lambda_N} \binom{N}{j}\|^2 \geq c_\alpha > 0, \quad \Lambda_N \subset V_N$

$F : \mathbb{R} \rightarrow \mathbb{R}, \quad \in I,$

$$(20) \quad \frac{1}{N} \sum_{x \in V_N} F \left( \binom{N}{\lambda + i\eta_0}(x) \right) \xrightarrow{N \rightarrow +\infty} \mathbb{E} \left( F \left( \frac{\text{Im } \mathcal{G}^{\lambda + i\eta_0}(o; o)}{\mathbb{E}(\text{Im } \mathcal{G}^{\lambda + i\eta_0}(o; o))} \right) \right) :$$

$\binom{N}{\lambda + i\eta_0}(x) \approx \binom{N}{\lambda + i\eta_0}(o; o), \quad \mathbb{E} \left( \frac{\text{Im } \mathcal{G}^{\lambda + i\eta_0}(o; o)}{\mathbb{E}(\text{Im } \mathcal{G}^{\lambda + i\eta_0}(o; o))} \right) = 1$

**Remark 14.** The results proven in actually hold for more general Schrödinger operators than adjacency matrices : one can consider weighted Laplacians (with conductances on the edges) and add a potential; in other words, on each  $G_N$ , we can consider a discrete Schrödinger operator  $\mathcal{H}_N$ . The limiting object in assumption (BSCT) is now a random



$$\langle a_N \rangle = \frac{1}{N} \sum_{x=1}^N a_N(x) \quad \text{and} \quad \langle |j^{(N)}(x)|^2 \rangle = \frac{1}{N} \sum_{x=1}^N |j^{(N)}(x)|^2$$

**Theorem 17** (ii). Let  $!$  be such that  $\sqrt{q} \geq (! + 1)2^{2\omega+45}$ .

(i) With probability  $\geq 1 - o(N^{-\omega+8})$  on the choice of the graph,

$$\|j\|_\infty \leq \frac{(\log N)^{121}}{\sqrt{N}}$$

for all eigenfunctions associated to eigenvalues such that  $|j \pm 2\sqrt{q}| > (\log N)^{-3/2}$ .

(ii) (Quantum Unique Ergodicity for random regular graphs) Given an observable  $a_N : \{1; \dots; N\} \rightarrow \mathbb{R}$ , we have, with probability  $\geq 1 - o(N^{-\omega+8})$  on the choice of the graph, for  $N$  large enough,

$$(21) \quad \left| \sum_{x=1}^N a_N(x) |j^{(N)}(x)|^2 - \langle a_N \rangle \right| \leq \frac{(\log N)^{250}}{N} \sqrt{\sum_x |a_N(x)|^2}$$

for all eigenfunctions associated to eigenvalues  $j \in (-2\sqrt{q} + ; 2\sqrt{q} - )$  (bulk eigenvalues).

In particular, if  $a_N = \mathbb{1}_{\Lambda_N}$  where  $\Lambda_N \subset \{1; \dots; N\}$ , we find

$$\left| \sum_{x \in \Lambda_N} |j^{(N)}(x)|^2 - \frac{|\Lambda_N|}{N} \right| \leq \frac{(\log N)^{250}}{N} \sqrt{|\Lambda_N|}$$

(ii)  $! > 8$   $q > 2^{128}$ .

$|\Lambda_N| > (\log N)^{500}$ .

$\infty$ ,

$$\left| \sum_{x \in \Lambda_N} |j^{(N)}(x)|^2 - \frac{|\Lambda_N|}{N} \right| \leq \frac{1}{\sqrt{N \log N}} \sqrt{|\Lambda_N|}$$

most  $|\Lambda_N| > N^{1/2}$ .

(1),  $a_N$   $\dots$  almost sure

**Remark 18.** Note that we emphasized Theorem 17 from 1A because our main concern here is the delocalization of eigenfunctions. The main focus of 1A is however on the universality of the local spectral statistics for random regular graphs. This would deserve a separate paper.

$G_N$   $N$  all  $j^{(N)}$



$\mathbb{S}^2$ ,  $g_1, \dots, g_k$  fi  $\Delta_{\mathbb{S}^2}$ ,  $SO(3)$ ,

$$T_k f(x) = \sum_{j=1}^k (f(g_j x) + f(g_j^{-1} x))$$

$\Delta_{\mathbb{S}^2}$ .

**Theorem 19** ( ). Assume that  $g_1, \dots, g_k$  generate a free subgroup of  $SO(3)$ .

For each  $\ell$ , let  $(\psi_j^{(\ell)})_{j=1}^{2\ell+1}$  be an orthonormal family of eigenfunctions of  $-\Delta_{\mathbb{S}^2}$  of eigenvalue  $\ell(\ell+1)$ , that are also eigenfunctions of  $T_k$ .

Then for any continuous function  $a$  on  $\mathbb{S}^2$ , we have

$$\frac{1}{2\ell+1} \sum_{j=1}^{2\ell+1} \left| \int_M a(x) |\psi_j^{(\ell)}(x)|^2 d\text{Vol}(x) - \int_M a(x) d\text{Vol}(x) \right|^2 \xrightarrow{\ell \rightarrow \infty} 0:$$

11.  $T_k$   $T_k$   $2k$   $N \sim 2$   $g_1, \dots, g_k$  1 1 .

**Remark 20.**  $T_k$  is not a pseudodifferential operator, so the argument sketched above to show that the basis  $(Y_\ell^m)_{\ell \geq 0, |m| \leq \ell}$  could not satisfy quantum ergodicity does not apply here.

**Remark 21.** We note that for very special choices of rotations – rotations that correspond to norm  $n$  elements in an order in a quaternion division algebra, the operators  $T_k$  are called Hecke operators. It has been conjectured by Böcherer, Sarnak, and Schulze-Pillot [2] that such joint eigenfunctions satisfy the much stronger quantum unique ergodicity property. This conjecture is still open.

3 . [ 2000,  $(n)$  infinite  $(\cdot, \cdot)$  almost-every ergodic component  $\Delta$  4 . [ 3 ,  $M$ , 10 one



$$\int_{S_N} |a(x)|^2 |e_i^{(N)}(x)|^2 dx - \langle a \rangle^2 = 0$$

4.3. Quantum ergodicity on Riemann surfaces of high genus. 11 12

Let  $(S_N)$  be a sequence of hyperbolic surfaces, whose genus (equivalently, volume) goes to  $\infty$ .  
 (EXP) Assume the first eigenvalue  $\lambda_1(N)$  of  $-\Delta$  on  $S_N$  is bounded away from 0 as  $N \rightarrow \infty$ .  
 (BSH) Assume there are few short geodesics; in other words,  $(S_N)$  converges in the Benjamini-Schramm sense to the hyperbolic disc: for any  $R > 0$ ,

$$\lim_{N \rightarrow +\infty} \frac{\text{Vol}\{x \in S_N : r(x) < R\}}{\text{Vol}(S_N)} = 0$$

where  $r(x)$  means the injectivity radius at  $x$ .

Fix an interval  $I \subset (0; +\infty)$ .

Let  $(e_i^{(N)})$  be an orthonormal basis of eigenfunctions of the Laplacian on  $S_N$ .

Let  $a = a_N : S_N \rightarrow \mathbb{C}$  be such that  $|a(x)| \leq 1$  for all  $x \in S_N$ . Then

$$\lim_{N \rightarrow +\infty} \frac{1}{\text{Vol}(S_N)} \sum_{\lambda_i(N) \in I} \left| \int_{S_N} a(x) |e_i^{(N)}(x)|^2 dx - \langle a \rangle \right|^2 = 0$$

where  $\langle a \rangle = \frac{1}{\text{Vol}(S_N)} \int_{S_N} a(x) dx$ .

(1; +\infty) L^2-  
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7 *Math. Z.* - 7 77

*Comm. Math. Phys.* 7 7

*Phys.* - *Comm. Math.*

77 *J. Phys. A* -

7 - 77 *Proc. Royal Soc. A*

ff

fi *Comm. Math. Phys.* - 7

*Comm. Math. Phys.* 7-

fi *Phys. Rev. Lett.* -

7 *Houches, 1989* 7- *Chaos et physique quantique (Les*

*Philosophical Magazine* -

*Comm. Math. Phys.* 7

*Doc. Math.* 7 - 7

*Acad. Sci. Paris* 7- *C. R. Math.*

*Spectral theory of infinite-area hyperbolic surfaces* *Progress in Math-*

*ematics* 7

*preprint* fi

7 *Comm. Math. Phys.* -

*C. R. A. S.* -

7 *preprint* 7 *l<sup>p</sup>*

7 *Int. Math. Res. Not.* 7

*Israel J. Math.* -

*Invent. Math.* - *Comm. Math. Phys.* 7-

*Ann. H. Poincaré, to appear*

*preprint*  
*Chaos et physique quantique (Les*

*Houches, 1989)* -7 *Ann. Appl.*

*Probab.* -

*Ann. Probab.* 7-

*Ann. Sci. Éc. Norm. Supér. (4)* 7 7 -

7 7 *preprint* 7

ff

*International Congress of Mathematicians. Vol. II* 7 - 7

ff

*Annalen der Physik* 7 *Verhandl. deut. physik. Ges.* 7

"

*Ann. Probab.* 7 - 7

*Comm. Math. Phys.* 7 -

*Ann. Probab.* 7 -

*Comm. Math. Phys.* -

*Comm. Math. Phys.* -

*J. Spectr. Theory* 7 - 7

7 *Duke Math. J.* 7 - 7 *Ann.*

*Physics* -

*Amer. J. Math.* - 7 *Nonlinearity*

- *Comm. Math. Phys.* -

7 *Ann. of Math. (2)*

7 - *Forum Math.* 7 -

*Zeitschrift f. Physik* 7 -

ff  
 - 7  
 Houches, 1989) - *Comm. Math. Phys.*  
 7  $L^p$  *Chaos et physique quantique (Les*  
 - *Adv. Math.*  
*Comm. Math.*  
 Phys. - 7  
 Fourier (Grenoble) 77-  $PSL_2(\mathbf{Z}) \backslash PSL_2(\mathbf{R})$  *Ann. Inst.*  
 7 *Amer. J. Math.* 7 -  
 7 *Comm. Math. Phys.* 7 - 7  
 7 *J. Mod. Dyn.* 7 -7  
 7 *Analysis on graphs and its applications*  
 77 *Proc. Sympos. Pure Math.* 7 -  
 7 *Comm. Math. Phys.* -  
 7 *Trans. Amer. Math. Soc.* -  
 7 *Adv. Math.* -  
 77 *Phys. Rev. Lett.* 7 7 -7 7  
 7 *Ann. Physics* 7 7 -  
 7 *Zeitschrift f. Physik*  
 -  
 7 *Duke Math. J.*  
 , ff  
*Ann. of Math. (2)* -  
 ff  
*Ann. of Math. (2)* - 7  
*Comm.*  
*Math. Phys.* 7 - 7  
*Ann. of Math.*  
 (2) - *Proceedings of*  
*the International Congress of Mathematicians. Volume I* - 7  
 7 *Combinatorica* - 77  
 - *Comm. Math. Phys.*  
 - 7 *Ann. of Math. (2)*  
*Appl.* - *Linear Algebra*

*Electron. J. Combin.*

*Proceedings of the International Congress of Mathematicians. Volume II* –

*J. Differential Geom.* 7–

– *Verhandl. deut. physik. Ges.*

*Duke Math. J.* 7 –

*Comm. Math. Phys.* –

*The Schur lectures (1992) (Tel Aviv)* Israel

*Math. Conf. Proc.* –

7 *Harmonic analysis and number theory*

*(Montreal, PQ, 1996)* *CMS Conf. Proc.*

7

*Bull. Amer. Math. Soc. (N.S.)* – 7

7 – 7 *Annalen der Physik (4)*

7 – 7 *Annalen der Physik (4)*

– –

*Annalen der Physik (4)* 7 7 –7

*Festschrift in honor of I. I. Piatetski-Shapiro on the occasion of his sixtieth birthday, Part II (Ramat Aviv, 1989)*

*Israel Math. Conf. Proc.* – 7

*Physica Scripta*

*J. Phys. A* 7 – 7

– *Séminaire Poincaré* –

7 *Uspehi Mat. Nauk* – 7

$L^p$

*J. Funct. Anal.* 77 –

*Duke Math. J.* 7– 7

*Iberoam.* 7 – *Rev. Mat.*


*Rev. Mat. Iberoam.* –

*Proceedings of the International Congress of Mathematicians. Volume II* 7–

$SL_2(\mathbb{Z}) \backslash \mathbb{H}$  *Ann. of Math. (2)* 7 –

*Random Structures Algorithms* –

$ff$   $L^\infty$  *Internat. Math. Res. Notices* 7 – 7 7

  
*Zeitschrift f. Physik* —  
*J.* — 7  
7  
*J. Funct. Anal.* 7 —  
  
*Phys.* 7 7 —  
  
*Duke Math.*  
  
*Comm. Math. Phys.* —7  
  
*Comm. Math.*

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